# SOME ALGEBRAIC PROPERTIES OF DUAL GENERALIZED QUATERNIONS ALGEBRA 

Hamid Mortazaasl, Mehdi Jafari and Yusuf Yayli*<br>Department of Mathematics<br>Islamic Azad University<br>Shabestar Branch, Shabestar, Iran<br>*Department of Mathematics<br>Faculty of Science<br>Ankara University<br>Ankara, Turkey


#### Abstract

Some algebraic properties of dual generalized quaternions are presented, and De-Moivre's and Euler's formulas for these quaternions are investigated. The solutions of equation $Q^{n}=1$ are discussed and it is shown that it has not solutions for a general unit dual generalized quaternion.


## 1. Introduction

Dual quaternions are powerful mathematical tools for the spatial analysis of rigid body motions. Dual numbers and dual quaternions were introduced © 2012 Pushpa Publishing House 2010 Mathematics Subject Classification: 30G35.
Keywords and phrases: De-Moivre's formula, dual generalized quaternion, Euler formula.
This work has been supported by a grant from the Islamic Azad University, Shabestar Branch with Number: 51953890703006.

Received April 19, 2012
in the 19th century by Clifford [2], as a tool for his geometrical investigation. Study [8] and Kotel'nikov [6] systematically applied the dual number and dual vector in their studies of line geometry and kinematics and independently discovered the transfer principle. The use of dual numbers, matrices and dual quaternions in instantaneous spatial kinematics are investigated in [9, 11]. Also, dual quaternion algebra is used to express pointline displacement operation [12].

The Euler's and De-Moivre's formulas for the complex numbers are generalized for quaternions in [1]. These formulas are also investigated for the cases of split and dual quaternions in [5, 7]. In [10], by the help of Euler's formula, circles in the plane are obtained and the sphere in 3-space is found by means of the exponential expansions. In this paper, after a review of some properties of dual generalized quaternions, De-Moivre's and Euler's formulas for these quaternions are studied. Solutions of the equation $Q^{n}=1$ are discussed and it is shown that there are no solutions for a general unit dual generalized quaternion. The relations between the powers of these quaternions are given.

## 2. Preliminaries

In this section, we give a brief summary of the generalized quaternions and dual generalized quaternions. For detailed information about these quaternions, we refer the reader to $[3,4]$.

Definition 2.1. A generalized quaternion $q$ is defined as

$$
q=a .+a_{1} i+a_{2} j+a_{3} k,
$$

where $a ., a_{1}, a_{2}$ and $a_{3}$ are real numbers and $1, i, j, k$ of $q$ may be interpreted as the four basic vectors of Cartesian set of coordinates; and they satisfy the non-commutative multiplication rules

$$
\begin{aligned}
& i^{2}=-\alpha, \quad j^{2}=-\beta, \quad k^{2}=-\alpha \beta, \\
& i j=k=-j i, \quad j k=\beta i=-k j
\end{aligned}
$$

and

$$
k i=\alpha j=-i k, \quad \alpha, \beta \in \mathbb{R} .
$$

The set of all generalized quaternions are denoted by $H_{\alpha \beta}$.
Definition 2.2. A dual generalized quaternion $Q$ is written as

$$
Q=A .1+A_{1} i+A_{2} j+A_{3} k,
$$

where $A ., A_{1}, A_{2}$ and $A_{3}$ are dual numbers and $i, j, k$ are quaternionic units which satisfy in the above equalities. As a consequence of this definition, a generalized dual quaternion $Q$ can also be written as:

$$
Q=q+\varepsilon q^{*}, \quad q, q^{*} \in H_{\alpha \beta},
$$

where $q$ and $q^{*}$, real and pure dual generalized quaternion components, respectively. A quaternion $Q=A .1+A_{1} i+A_{2} j+A_{3} k$ is pieced into two parts with scalar piece $S_{Q}=A$. and vectorial piece $\vec{V}_{Q}=A_{1} i+A_{2} j+A_{3} k$. We also write $Q=S_{Q}+\vec{V}_{Q}$. The conjugate of $Q=S_{Q}+\vec{V}_{Q}$ is then defined as $\bar{Q}=S_{Q}-\vec{V}_{Q}$. If $S_{Q}=0$, then $Q$ is called pure dual generalized quaternion, we may be called its dual generalized vector.

Dual quaternionic multiplication of two dual quaternions $Q=S_{Q}+\vec{V}_{Q}$ and $P=S_{P}+\vec{V}_{P}$ is defined:

$$
\begin{aligned}
Q P= & S_{Q} S_{P}-g\left(\vec{V}_{Q}, \vec{V}_{P}\right)+S_{P} \vec{V}_{Q}+S_{Q} \vec{V}_{P}+\vec{V}_{Q} \wedge \vec{V}_{P} \\
= & \text { A.B. }-\left(\alpha A_{1} B_{1}+\beta A_{2} B_{2}+\alpha \beta A_{3} B_{3}\right)+A \cdot\left(B_{1}, B_{2}, B_{3}\right) \\
& +B \cdot\left(A_{1}, A_{2}, A_{3}\right) \\
& +\left(\beta\left(A_{2} B_{3}-A_{3} B_{2}\right), \alpha\left(A_{3} B_{1}-A_{1} B_{3}\right),\left(A_{1} B_{2}-A_{2} B_{1}\right)\right) .
\end{aligned}
$$

The norm of $Q$ is defined as $N_{Q}=Q \bar{Q}=\bar{Q} Q=A^{2}+\alpha A_{1}^{2}+\beta A_{2}^{2}$ $+\alpha \beta A_{3}^{2}$. If $N_{Q}=1$, then $Q$ is called a unit generalized dual quaternion. The set of all dual generalized quaternions (DGQ) is denoted by $\tilde{H}_{\alpha \beta}$.

## 3. De Moiver's Formula for DGQ

We investigate the properties of the dual generalized quaternions in two different cases:

Case 1. Let $\alpha, \beta$ be positive numbers.
Definition 3.1. Let $\hat{S}_{D}^{3}$ be the set of all unit dual generalized quaternions and $\hat{S}_{D}^{2}$ be the set of unit dual generalized vector, that is,

$$
\begin{aligned}
& \hat{S}_{D}^{3}=\left\{Q \in \tilde{H}_{\alpha \beta}: N_{Q}=1\right\} \subset \tilde{H}_{\alpha \beta} \\
& \hat{S}_{D}^{2}=\left\{\vec{V}_{Q}=\left(A_{1}, A_{2}, A_{3}\right): g\left(\vec{V}_{Q}, \vec{V}_{Q}\right)=\alpha A_{1}^{2}+\beta A_{2}^{2}+\alpha \beta A_{3}^{2}=1\right\} .
\end{aligned}
$$

Definition 3.2. Every nonzero unit dual generalized quaternion can be written in the polar form

$$
\begin{aligned}
Q & =A_{1}+A_{1} i+A_{2} j+A_{3} k \\
& =\cos \phi+\vec{W} \sin \phi,
\end{aligned}
$$

where $\cos \phi=A, \sin \phi=\sqrt{\alpha A_{1}^{2}+\beta A_{2}^{2}+\alpha \beta A_{3}^{2}} \cdot \phi=\varphi+\varepsilon \varphi^{*}$ is a dual angle and the unit dual generalized vector $\vec{W}$ is given by

$$
\vec{W}=\frac{A_{1} i+A_{2} j+A_{3} j}{\sqrt{\alpha A_{1}^{2}+\beta A_{2}^{2}+\alpha \beta A_{3}^{2}}}=\frac{A_{1} i+A_{2} j+A_{3} j}{\sqrt{1-A^{2}}},
$$

with $\alpha A_{1}^{2}+\beta A_{2}^{2}+\alpha \beta A_{3}^{2} \neq 0$.
For any $\vec{W} \in \hat{S}_{D}^{2}$, since $g(\vec{W}, \vec{W})=1$ and $\vec{W} \wedge \vec{W}=0, \quad \vec{W}^{2}=-1$. Therefore, any $\vec{W} \in \hat{S}_{D}^{2}$ is of order 4, i.e., $\vec{W}^{4}=1$. We have a natural generalization of Euler's formula for dual generalized quaternions

$$
e^{\vec{W} \phi}=1+\vec{W} \phi-\frac{\phi^{2}}{2!}-\vec{W} \frac{\phi^{3}}{3!}+\frac{\phi^{4}}{4!}+\cdots
$$

$$
\begin{aligned}
& =\left(1-\frac{\phi^{2}}{2!}+\frac{\phi^{4}}{4}-\cdots\right)+\vec{W}\left(\phi-\frac{\phi^{3}}{3}+\frac{\phi^{5}}{5!}-\cdots\right) \\
& =\cos \phi+\vec{W} \sin \phi,
\end{aligned}
$$

for any dual number $\phi$. For detailed information about Euler's formula, see [10].

Note that $\vec{W}$ is a unit dual generalized vector to which a directed line in $E_{\alpha \beta}^{3}$ corresponds by means of the generalized Study map.

Lemma 3.1. For any $\vec{W} \in \hat{S}_{D}^{2}$, we have

$$
(\cos \phi+\vec{W} \sin \phi)(\cos \psi+\vec{W} \sin \psi)=\cos (\phi+\psi)+\vec{W} \sin (\phi+\psi) .
$$

Proof. See [5] for a similar proof.
Theorem 3.2 (De-Moivre's formula). Let $Q=e^{\vec{W} \phi}=\cos \phi+\vec{W} \sin \phi$ $\in \hat{S}_{D}^{3}$, where $\phi=\varphi+\varepsilon \varphi^{*}$ is dual angle and $\vec{W} \in \hat{S}_{D}^{2}$. Then for every integer n,

$$
Q^{n}=\cos n \phi+\vec{W} \sin n \phi .
$$

Proof. We use induction on positive integers $n$. Assume that $Q^{n}=\cos n \phi$ $+\vec{W} \sin n \phi$ holds. Then

$$
\begin{aligned}
Q^{n+1} & =(\cos \phi+\vec{W} \sin \phi)^{n}(\cos \phi+\vec{W} \sin \phi) \\
& =(\cos n \phi+\vec{W} \sin n \phi)(\cos \phi+\vec{W} \sin \phi) \\
& =\cos (n \phi+\phi)+\vec{W} \sin (n \phi+\phi) \\
& =\cos (n+1) \phi+\vec{W} \sin (n+1) \phi .
\end{aligned}
$$

Hence, the formula is true. Moreover, since

$$
\begin{aligned}
Q^{-1} & =\cos \phi-\vec{W} \sin \phi \\
Q^{-n} & =\cos (-n \phi)+\vec{W} \sin (-n \phi) \\
& =\cos n \phi-\vec{W} \sin n \phi
\end{aligned}
$$

the formula holds for all integers.
Special case. If $\alpha=\beta=1$, then Theorem 3.2 holds for dual quaternions (see [5]).

Every generalized dual quaternion is separated into two cases:
(1) Dual generalized quaternion with dual angles $\left(\phi=\varphi+\varepsilon \varphi^{*}\right)$; i.e.,

$$
Q=\sqrt{N_{Q}}(\cos \phi+\vec{W} \sin \phi) .
$$

(2) Dual generalized quaternions with real angles $\left(\phi=\varphi, \varphi^{*}=0\right)$; i.e.,

$$
Q=\sqrt{N_{Q}}(\cos \varphi+\vec{W} \sin \varphi) .
$$

Theorem 3.3. Let $Q=\cos \varphi+\vec{W} \sin \varphi \in \hat{S}_{D}^{3}$. De-Moivre's formula implies that there are uncountably many unit dual generalized quaternions satisfying $Q^{n}=1$ for $n \geq 3$.

Proof. For every $\vec{W} \in \hat{S}_{D}^{2}$, the unit dual generalized quaternion

$$
Q=\cos \frac{2 \pi}{n}+\vec{W} \sin \frac{2 \pi}{n}
$$

is of order $n$. For $n=1$ or $n=2$, the dual generalized quaternion $Q$ is independent of $\vec{W}$.

## Example 3.1.

$$
Q=\frac{1}{\sqrt{2}}+\frac{1}{2}\left(\frac{1}{\sqrt{\alpha}}, \frac{1}{\sqrt{\beta}}, \varepsilon\right)=\cos \frac{\pi}{4}+\frac{1}{\sqrt{2}}\left[\frac{1}{\sqrt{\alpha}} i+\frac{1}{\sqrt{\beta}} j+\varepsilon k\right] \sin \frac{\pi}{4}
$$

is of order 8 and

$$
Q_{2}=-\frac{1}{2}+\left(2 \varepsilon, \frac{1}{\sqrt{2 \beta}}, \frac{1}{2 \sqrt{\alpha \beta}}\right)=\cos \frac{2 \pi}{3}+\frac{1}{\sqrt{3}}\left[2 \varepsilon i+\sqrt{\frac{2}{\beta}} j+\frac{1}{\sqrt{\alpha \beta}} k\right] \sin \frac{2 \pi}{3}
$$

is of order 3.
Also, we find the $n$th root of $Q=\cos \varphi+\vec{W} \sin \varphi \in \hat{S}_{D}^{3}$. The equation $X^{n}=Q$ has $n$ roots. Thus,

$$
X=\cos \left(\frac{\varphi+2 k \pi}{n}\right)+\vec{W} \sin \left(\frac{\varphi+2 k \pi}{n}\right), \quad k=0, \ldots, n-1 .
$$

Corollary 3.4. The equation $Q^{n}=1$ does not have any solution for $a$ general unit dual generalized quaternions.

Example 3.2. Let $Q=\cos \left(\frac{\pi}{3}+\varepsilon\right)+\vec{W} \sin \left(\frac{\pi}{3}+\varepsilon\right)$ be a unit dual generalized quaternion. There is no $n(n>0)$ such that $Q^{n}=1$.

Theorem 3.5. Let $Q$ be a unit dual generalized quaternion with the polar form $Q=\cos \varphi+\vec{W} \sin \varphi$. If $m=\frac{2 \pi}{\varphi} \in \mathbb{Z}^{+}-\{1\}$, then $n \equiv p(\bmod m)$ is possible if and only if $Q^{n}=Q^{p}$.

Proof. Let $n \equiv p(\bmod m)$. Then we have $n=a . m+p$, where $a \in \mathbb{Z}$,

$$
\begin{aligned}
Q^{n} & =\cos n \varphi+\vec{W} \sin n \varphi \\
& =\cos (a m+p) \varphi+\vec{W} \sin (a m+p) \varphi \\
& =\cos \left(a \frac{2 \pi}{\varphi}+p\right) \varphi+\vec{W} \sin \left(a \frac{2 \pi}{\varphi}+p\right) \varphi \\
& =\cos (p \varphi+a 2 \pi)+\vec{W} \sin (p \varphi+a 2 \pi) \\
& =\cos (p \varphi)+\vec{W} \sin (p \varphi) \\
& =Q^{p} .
\end{aligned}
$$

Now suppose $Q^{n}=\cos n \varphi+\vec{W} \sin n \varphi$ and $Q^{p}=\cos p \varphi+\vec{W} \sin p \varphi$. Since $Q^{n}=Q^{p}$, we have $\cos n \varphi=\cos p \varphi$ and $\sin n \varphi=\sin p \varphi$, which means $n \varphi=p \varphi+2 \pi a, a \in \mathbb{Z}$. Thus, $n=a \frac{2 \pi}{\varphi}+p, n \equiv p(\bmod m)$.

Example 3.3. Letting $Q=\frac{1}{\sqrt{2}}+\frac{1}{2}\left(\frac{1}{\sqrt{\alpha}}, \frac{1}{\sqrt{\beta}}, \varepsilon\right) \in \hat{S}_{D}^{3}$ from Theorem 3.5, $m=\frac{2 \pi}{\pi / 4}=8$, we have

$$
\begin{aligned}
& Q=Q^{9}=Q^{17}=\cdots \\
& Q^{2}=Q^{10}=Q^{18}=\cdots \\
& Q^{3}=Q^{11}=Q^{19}=\cdots \\
& Q^{4}=Q^{12}=Q^{20}=\cdots=-1 \\
& \cdots \\
& Q^{8}=Q^{16}=Q^{24}=\cdots=1 .
\end{aligned}
$$

Case 2. Let $\alpha$ be a positive number and $\beta$ be a negative number.
In this case, for a dual generalized quaternion $Q=A+A_{1} i+A_{2} j+A_{3} k$, we can consider three different cases:

Case A. The norm of dual generalized quaternion is negative, i.e.,

$$
N_{Q}=A^{2}+\alpha A_{1}^{2}+\beta A_{2}^{2}+\alpha \beta A_{3}^{2}<0,
$$

since $0<A^{2}<-\alpha A_{1}^{2}-\beta A_{2}^{2}-\alpha \beta A_{3}^{2}$ thus $\alpha A_{1}^{2}+\beta A_{2}^{2}+\alpha \beta A_{3}^{2}<0$. In this case, the polar form of $Q$ is defined as

$$
Q=r(\sinh \Psi+\vec{W} \cosh \Psi)
$$

where we assume

$$
\begin{aligned}
& r=\sqrt{\left|N_{Q}\right|}=\sqrt{\left|A^{2}+\alpha A_{1}^{2}+\beta A_{2}^{2}+\alpha \beta A_{3}^{2}\right|}, \\
& \sinh \Psi=\frac{A}{\sqrt{\mid N_{Q}}}, \cosh \Psi=\frac{\sqrt{-\alpha A_{1}^{2}-\beta A_{2}^{2}-\alpha \beta A_{3}^{2}}}{\sqrt{\mid N_{Q}}}
\end{aligned}
$$

The unit dual vector $\vec{W}$ (axis of quaternion) is defined as

$$
\vec{W}=\left(w_{1}, w_{2}, w_{3}\right)=\frac{1}{\sqrt{-\alpha A_{1}^{2}-\beta A_{2}^{2}-\alpha \beta A_{3}^{2}}}\left(A_{1}, A_{2}, A_{3}\right) .
$$

Special case. If $\alpha=1, \beta=-1$ and the norm be negative number, i.e.,

$$
N_{Q}=A^{2}+A_{1}^{2}-A_{2}^{2}-A_{3}^{2}<0,
$$

then the quaternion $Q$ is a dual split quaternion with negative norm and its polar forms

$$
Q=\sqrt{\left|N_{Q}\right|}(\sinh \Psi+\vec{W} \cosh \Psi) .
$$

Theorem 3.6 (De-Moivre's formula). Let $Q=r(\sinh \Psi+\vec{W} \cosh \Psi)$ be a dual generalized quaternion with $N_{Q}<0$. Then for every integer $n$,

$$
Q^{n}=r^{n}(\sinh n \Psi+\vec{W} \cosh n \Psi)
$$

Proof. The proof follows immediately from the induction.
Case B. The norm of dual generalized quaternion is positive and the norm of its vector part to be negative, i.e.,

$$
N_{Q}>0, \quad \vec{V}_{Q}=\alpha A_{1}^{2}+\beta A_{2}^{2}+\alpha \beta A_{3}^{2}<0
$$

In this case, the polar form of $Q$ is defined as

$$
Q=r(\cosh \Phi+\vec{W} \sinh \Phi),
$$

where we assume

$$
\begin{aligned}
& r=\sqrt{N_{Q}}=\sqrt{A^{2}+\alpha A_{1}^{2}+\beta A_{2}^{2}+\alpha \beta A_{3}^{2}}, \\
& \cosh \Phi=\frac{A .}{\sqrt{N_{Q}}}, \sinh \Phi=\frac{\sqrt{-\alpha A_{1}^{2}-\beta A_{2}^{2}-\alpha \beta A_{3}^{2}}}{\sqrt{N_{Q}}} .
\end{aligned}
$$

The unit dual vector $\vec{W}$ (axis of quaternion) is defined as

$$
\vec{W}=\left(w_{1}, w_{2}, w_{3}\right)=\frac{1}{\sqrt{-\alpha A_{1}^{2}-\beta A_{2}^{2}-\alpha \beta A_{3}^{2}}}\left(A_{1}, A_{2}, A_{3}\right) .
$$

Theorem 3.7. Let $Q=r(\cosh \Phi+\vec{W} \sinh \Phi)$ be a dual generalized quaternion with $N_{Q}>0$ and $\vec{V}_{Q}<0$. Then for every integer $n$,

$$
Q^{n}=r^{n}(\cosh n \Phi+\vec{W} \sinh n \Phi)
$$

Proof. We use induction on positive integers n. Assume that $Q^{n}=r^{n}(\cosh n \Phi+\vec{W} \sinh n \Phi)$ holds. Then

$$
\begin{aligned}
Q^{n+1}= & r^{n}(\cosh n \Phi+\vec{W} \sinh n \Phi) \cdot r(\cosh \Phi+\vec{W} \sinh \Phi) \\
= & r^{n+1}(\cosh n \Phi+\vec{W} \sinh n \Phi)(\cosh \Phi+\vec{W} \sinh \Phi) \\
= & r^{n+1}(\cosh n \Phi \cosh \Phi+\sinh n \Phi \sinh \Phi \\
& +\vec{W}(\cosh n \Phi \sinh \Phi+\sinh n \Phi \cosh \Phi) \\
= & r^{n+1}(\cosh (n+1) \Phi+\vec{W} \sinh (n+1) \Phi) .
\end{aligned}
$$

Hence, the formula is true.
Case C. The norm of dual generalized quaternion is positive and the norm of its vector part to be positive, i.e.

$$
N_{Q}>0, \quad \vec{V}_{Q}=\alpha A_{1}^{2}+\beta A_{2}^{2}+\alpha \beta A_{3}^{2}>0
$$

In this case, the polar form of $Q$ is defined as

$$
Q=r(\cos \Theta+\vec{W} \sin \Theta)
$$

where we assume

$$
\begin{aligned}
& r=\sqrt{N_{Q}}=\sqrt{A^{2}+\alpha A_{1}^{2}+\beta A_{2}^{2}+\alpha \beta A_{3}^{2}}, \\
& \cos \Theta=\frac{A .}{\sqrt{N_{Q}}}, \quad \sin \Theta=\frac{\sqrt{\alpha A_{1}^{2}+\beta A_{2}^{2}+\alpha \beta A_{3}^{2}}}{\sqrt{N_{Q}}} .
\end{aligned}
$$

The unit dual vector $\vec{W}$ (axis of quaternion) is defined as

$$
\vec{W}=\left(w_{1}, w_{2}, w_{3}\right)=\frac{1}{\sqrt{\alpha A_{1}^{2}+\beta A_{2}^{2}+\alpha \beta A_{3}^{2}}}\left(A_{1}, A_{2}, A_{3}\right) .
$$

Theorem 3.8. Let $Q=r(\cos \Theta+\vec{W} \sin \Phi)$ be a dual generalized quaternion with $N_{Q}>0$ and $\vec{V}_{Q}>0$. Then for every integer $n$,

$$
Q^{n}=r^{n}(\cos n \Theta+\vec{W} \sin n \Theta) .
$$

Proof. The proof follows immediately from the induction.

## Conclusion

In this paper, we defined and gave some of algebraic properties of dual generalized quaternion and investigated the Euler's and De-Moivre's formulas for these quaternions in several cases. The relation between the powers of DGQ is given in Theorem 3.5. We also showed that the equation $Q^{n}=1$ does not have solution for a general unit dual generalized quaternion (Corollary 3.4).

## References

[1] E. Cho, De-Moivre formula for quaternions, Appl. Math. Lett. 11(6) (1998), 33-35.
[2] W. Clifford, Preliminary sketch of biquaternions, Proc. London Math. Soc. 10 (1873), 381-395.
[3] M. Jafari and Y. Yayli, Rotation in four dimensions via generalized Hamilton operators, Accepted for Publication in Kuwait J. of Sci. and Eng. 40(1A) (2013), 35-47.
[4] M. Jafari and Y. Yayli, Dual generalized quaternions in spatial kinematics. 41st Annual Iranian Math. Conference, Urmia, Iran, 12-15 Sep. 2010.
[5] H. Kabadayi and Y Yayli, De-Moivre's formula for dual quaternions, Kuwait J. Sci. and Tech. 38(1) (2011), 15-23.
[6] A. P. Kotel’nikov, Vintovoe Schislenie i Nikotoriya Prilozheniye evo k geometrie i mechaniki, Kazan, 1895.
[7] M. Ozdemir, The roots of a split quaternion, Appl. Math. Lett. 22 (2009), 258-263.
[8] E. Study, Von Den Bewegungen und Umlegungen, Mathematische Annalen 39 (1891), 441-564.
[9] G. R. Veldkamp, On the use of dual numbers, vectors and matrices in instantaneous spatial kinematics, Mech. Mach. Theory 11 (1976), 141-156.
[10] J. Whittlesey and K. Whittlesey, Some geometrical generalizations of Euler's formula, Int. J. Math. Edu. Sci. Tech. 21(3) (1990), 461-468.
[11] A. T. Yang and F. Freudenstein, Application of dual-number quaternion algebra to the analysis of spatial mechanisms, ASME J. Appl. Mech. 86E (1964), 300-308.
[12] Y. Zhang and K. Ting, On point-line geometry and displacement, Mech. Mach. Theory 39 (2004), 1033-1050.

