



AN INTEGRAL FORMULA OF THE MELLIN TRANSFORM TYPE INVOLVING THE EXTENDED WRIGHT-BESSEL FUNCTION

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Received: August 30, 2017; Accepted: October 14, 2017

2010 Mathematics Subject Classification: 33C45, 33C60, 33E12.

Keywords and phrases: Wright-Bessel function, extended Wright-Bessel functions, Mittag-Leffler function and its extensions, generalized hypergeometric function, generalized (Wright) hypergeometric function, Gamma function, Pochhammer symbol, Mellin transform type integral formula.

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Abstract

We present a new Mellin transform type integral formula involving the generalized Wright-Bessel (or Bessel-Maitland) function defined by Ghayasuddin et al. [7, 8] expressed in terms of the generalized (Wright) hypergeometric function. The integral formula presented here, being very general, can be specialized to yield numerous integral formulas involving extended Wright-Bessel function and extended Mittag-Leffler functions, only three of which are demonstrated.

1. Introduction and Preliminaries

Recently, numerous integral formulas involving a variety of special functions and their generalizations have been presented (see, e.g., [2, 7-9, 11]). Also, many integral formulas associated with the Bessel functions of several kinds and their extensions have been established (see, e.g., [3-5]). Those integrals involving Wright-Bessel functions play important roles in many branches of theoretical and applied physics and engineering. Very recently, Abouzaid et al. [1] and Khan and Kashmin [10] have presented certain interesting and new classes of integral formulas involving the generalized Wright-Bessel function, which are expressed in terms of the generalized (Wright) hypergeometric function. In the present sequel to the aforementioned investigations, we present a Mellin transform type integral formula involving generalized Wright-Bessel functions and its variant, which are expressed in terms of the generalized (Wright) hypergeometric function ${}_p\Psi_q$ and the generalized hypergeometric function ${}_pF_q$, respectively. Some particular cases of our main results are also considered.

The Wright-Bessel function $J_v^\mu(z)$ is defined by the following series (see [12, equation (8.3)]):

$$J_v^\mu(z) = \sum_{n=0}^{\infty} \frac{(-z)^n}{n! \Gamma(\mu n + v + 1)}. \quad (1.1)$$

Singh et al. [19] introduced the following generalization of Wright-Bessel function:

$$J_{\nu, q}^{\mu, \gamma} = \sum_{n=0}^{\infty} \frac{(\gamma)_{qn} (-z)^n}{\Gamma(\mu n + \nu + 1) n!}$$

$$(\Re(\mu) \geq 0, \Re(\nu) \geq -1, \Re(\gamma) \geq 0, q \in (0, 1) \cup \mathbb{N}), \quad (1.2)$$

where $(\lambda)_{\nu}$ is the Pochhammer symbol defined (for $\lambda, \nu \in \mathbb{C}$) by (see [21, p. 2 and p. 5]):

$$(\lambda)_{\nu} := \frac{\Gamma(\lambda + \nu)}{\Gamma(\lambda)} \quad (\lambda + \nu \in \mathbb{C} \setminus \mathbb{Z}_0^-)$$

$$= \begin{cases} 1 & (\nu = 0) \\ \lambda(\lambda + 1) \cdots (\lambda + n - 1) & (\nu = n \in \mathbb{N}) \end{cases} \quad (1.3)$$

and $\Gamma(\lambda)$ is the familiar Gamma function. Here and in the following, let \mathbb{C} , \mathbb{R}^+ , \mathbb{N} and \mathbb{Z}_0^- be the sets of complex numbers, positive real numbers, positive integers and non-positive integers, respectively, and let $\mathbb{N}_0 := \mathbb{N} \cup \{0\}$.

In the sequel of the above-cited works, Ghayasuddin et al. [7] introduced and investigated a new extension of Wright-Bessel function as follows:

$$J_{\nu, \gamma, \delta}^{\mu, q, p}(z) = \sum_{n=0}^{\infty} \frac{(\gamma)_{qn} (-z)^n}{\Gamma(\mu n + \nu + 1) (\delta)_{pn}}$$

$$(\Re(\mu) \geq 0, \Re(\nu) \geq -1, \Re(\gamma) \geq 0, \Re(\delta) \geq 0,$$

$$p, q \in \mathbb{R}^+, q < \Re(\mu) + p). \quad (1.4)$$

We consider some special cases of the extended Wright-Bessel function $J_{\nu, \gamma, \delta}^{\mu, q, p}(z)$ in (1.4):

(i) We have

$$J_{v-1, q, p}^{\mu, \gamma, \delta}(-z) = E_{v, q, p}^{\mu, \gamma, \delta}(z), \quad (1.5)$$

where $E_{v, q, p}^{\mu, \gamma, \delta}(z)$ is the extended Mittag-Leffler function in [18].

(ii) We get

$$J_{v-1, q, 1}^{\mu, \gamma, 1}(-z) = E_{v, q}^{\mu, \gamma}(z), \quad (1.6)$$

where $E_{v, q}^{\mu, \gamma}(z)$ is the extended Mittag-Leffler function in [20].

(iii) We find

$$J_{v-1, 1, 1}^{\mu, \gamma, 1}(-z) = E_{\mu, v}^{\gamma}(z), \quad (1.7)$$

where $E_{\mu, v}^{\gamma}(z)$ is the extended Mittag-Leffler function in [15].

(iv) We obtain

$$J_{v-1, 1, 1}^{\mu, 1, 1}(-z) = E_{\mu, v}(z), \quad (1.8)$$

where $E_{\mu, v}(z)$ is the Mittag-Leffler function in [23].

(v) We see

$$J_{0, 1, 1}^{\mu, 1, 1}(-z) = E_{\mu}(z), \quad (1.9)$$

where $E_{\mu}(z)$ is the Mittag-Leffler function in [13].

An interesting generalization of the generalized hypergeometric series ${}_pF_q$ (see, e.g., [21, Section 1.5]) is due to Fox [6] and Wright [24-26] who studied the asymptotic expansion of the generalized (Wright) hypergeometric function defined by (see [22, p. 21]; see also [17])

$${}_p\Psi_q\left[\begin{matrix}(\alpha_1, A_1), \dots, (\alpha_p, A_p); \\ (\beta_1, B_1), \dots, (\beta_q, B_q); \end{matrix} z\right] = \sum_{k=0}^{\infty} \frac{\prod_{j=1}^p \Gamma(\alpha_j + A_j k)}{\prod_{j=1}^q \Gamma(\beta_j + B_j k)} \frac{z^k}{k!}, \quad (1.10)$$

where the coefficients $A_1, \dots, A_p \in \mathbb{R}^+$ and $B_1, \dots, B_q \in \mathbb{R}^+$ such that

$$1 + \sum_{j=1}^q B_j - \sum_{j=1}^p A_j \geq 0. \quad (1.11)$$

A special case of (1.10) is

$${}_p\Psi_q\left[\begin{matrix}(\alpha_1, 1), \dots, (\alpha_p, 1); \\ (\beta_1, 1), \dots, (\beta_q, 1); \end{matrix} z\right] = \frac{\prod_{j=1}^p \Gamma(\alpha_j)}{\prod_{j=1}^q \Gamma(\beta_j)} {}_pF_q\left[\begin{matrix}\alpha_1, \dots, \alpha_p; \\ \beta_1, \dots, \beta_q; \end{matrix} z\right]. \quad (1.12)$$

We also need to recall the following integral formula of the Mellin transform type (see [14, p. 22, Entry 2.47]):

$$\int_0^{\infty} x^{\mu-1} (x+a+\sqrt{x^2+2ax})^{-\lambda} dx = 2\lambda a^{-\lambda} \left(\frac{a}{2}\right)^{\mu} \frac{\Gamma(2\mu)\Gamma(\lambda-\mu)}{\Gamma(1+\lambda+\mu)} \\ (a \in \mathbb{R}^+; 0 < \Re(\mu) < \lambda). \quad (1.13)$$

2. Integral Formulas

Here, we establish a Mellin transform type integral formula involving the extended Wright-Bessel function (1.4) and its variant.

Theorem 2.1. *Let $\delta, \eta, v \in \mathbb{C}$ with $\Re(\delta) > 0$, $\Re(\eta) > 0$. Also, let $a, p, q, \gamma, \lambda, \mu \in \mathbb{R}^+$ with $\mu + p - q \geq 0$, $0 < \Re(\delta) < \lambda$, and $\Re(v) > -1 - \mu$. Then*

$$\begin{aligned}
& \int_0^\infty x^{\delta-1} (x+a+\sqrt{x^2+2ax})^{-\lambda} J_{\nu, \eta, p}^{\mu, \gamma, q} \left(\frac{y}{x+a+\sqrt{x^2+2ax}} \right) dx \\
&= 2^{1-\delta} a^{\delta-\lambda} \frac{\Gamma(2\delta)\Gamma(\eta)}{\Gamma(\gamma)} \\
&\quad \times {}_4\Psi_4 \left[\begin{matrix} (\gamma, q), (\lambda-\delta, 1), (\lambda+1, 1), (1, 1); \\ (\nu+1, \mu), (\lambda, 1), (1+\lambda+\delta, 1), (\eta, p); \end{matrix} -\frac{y}{a} \right]. \quad (2.1)
\end{aligned}$$

Proof. Let \mathcal{L} be the left side of (2.1). Using (1.4) and changing the order of integral and summation, which is verified under the given conditions, we obtain

$$\mathcal{L} = \sum_{n=0}^{\infty} \frac{(\gamma)_{qn} (-y)^n}{\Gamma(\mu n + \nu + 1) (\eta)_{pn}} \int_0^\infty x^{\delta-1} (x+a+\sqrt{x^2+2ax})^{-\lambda-n} dx. \quad (2.2)$$

Applying the formula (1.13) to the integral in (2.2), we obtain

$$\begin{aligned}
\mathcal{L} &= 2^{1-\delta} a^{\delta-\lambda} \frac{\Gamma(2\delta)\Gamma(\eta)}{\Gamma(\gamma)} \\
&\quad \times \sum_{n=0}^{\infty} \frac{\Gamma(\gamma+qn)\Gamma(\lambda+n-\delta)\Gamma(\lambda+n+1)\Gamma(n+1)}{\Gamma(\mu n + \nu + 1)\Gamma(\lambda+n)\Gamma(1+\lambda+n+\delta)\Gamma(\eta+pn)n!} \left(-\frac{y}{a} \right)^n,
\end{aligned}$$

which, upon expressing in terms of the generalized (Wright) hypergeometric function (1.10), leads to the right side of (2.1). \square

Under a little stronger condition, by using the following multiplication formula (see, e.g., [21, p. 6, equation (30)]):

$$(\lambda)_{mn} = m^{mn} \prod_{j=1}^m \left(\frac{\lambda+j-1}{m} \right)_n \quad (m \in \mathbb{N}; n \in \mathbb{N}_0),$$

the integral formula (2.1) can be expressed in terms of the generalized hypergeometric function ${}_pF_q$ which is asserted in the following theorem.

Theorem 2.2. Let $\delta, \eta, \nu \in \mathbb{C}$ with $\Re(\delta) > 0$, $\Re(\eta) > 0$. Also, let $a, \gamma, \lambda \in \mathbb{R}^+$ and $p, q, \mu \in \mathbb{N}$ with $\mu + p - q \geq 0$, $0 < \Re(\delta) < \lambda$, and $\Re(\nu) > -1 - \mu$. Then

$$\begin{aligned} & \int_0^\infty x^{\delta-1} (x + a + \sqrt{x^2 + 2ax})^{-\lambda} J_{\nu, \eta, p}^{\mu, \gamma, q} \left(\frac{y}{x + a + \sqrt{x^2 + 2ax}} \right) dx \\ &= 2^{1-\delta} a^{\delta-\lambda} \frac{\Gamma(\lambda+1)\Gamma(2\delta)\Gamma(\lambda-\delta)}{\Gamma(\lambda)\Gamma(\nu+1)\Gamma(\lambda+\delta+1)} \\ & \times {}_{q+3}F_{\mu+p+2} \left[\begin{matrix} \Delta(q; \gamma), \lambda+1, \lambda-\delta, 1; \\ \Delta(\mu; \nu+1), \Delta(p; \eta), \lambda, 1+\lambda+\delta; \end{matrix} \frac{q^q y}{\mu^\mu p^p a} \right], \quad (2.3) \end{aligned}$$

where $\Delta(m; \ell)$ abbreviates the array of m parameters

$$\frac{\ell}{m}, \frac{\ell+m}{m}, \dots, \frac{\ell+m-1}{m} \quad (m \in \mathbb{N}).$$

3. Special Cases

Here, among numerous special cases of the results in Section 2, only three of which are presented.

(i) Using the relation (1.5) in the formula (2.1), we have

$$\begin{aligned} & \int_0^\infty x^{\delta-1} (x + a + \sqrt{x^2 + 2ax})^{-\lambda} E_{\nu, \eta, p}^{\mu, \gamma, q} \left(\frac{y}{x + a + \sqrt{x^2 + 2ax}} \right) dx \\ &= 2^{1-\delta} a^{\delta-\lambda} \frac{\Gamma(2\delta)\Gamma(\eta)}{\Gamma(\gamma)} \\ & \times {}_4\Psi_4 \left[\begin{matrix} (\gamma, q), (\lambda-\delta, 1), (\lambda+1, 1), (1, 1); \\ (\nu, \mu), (\lambda, 1), (1+\lambda+\delta, 1), (\eta, p); \end{matrix} \frac{y}{a} \right], \quad (3.1) \end{aligned}$$

provided $\delta, \eta, \nu \in \mathbb{C}$ with $\min\{\Re(\delta), \Re(\eta)\} > 0$ and $a, p, q, \gamma, \lambda, \mu \in \mathbb{R}^+$ with $\mu + p - q \geq 0$, $0 < \Re(\delta) < \lambda$, and $\Re(\nu) > -\mu$.

(ii) Using the relation (1.5) in the formula (2.3), we obtain

$$\begin{aligned} & \int_0^\infty x^{\delta-1} (x+a+\sqrt{x^2+2ax})^{-\lambda} E_{\nu, \eta, p}^{\mu, \gamma, q} \left(\frac{y}{x+a+\sqrt{x^2+2ax}} \right) dx \\ &= 2^{1-\delta} a^{\delta-\lambda} \frac{\Gamma(\lambda+1)\Gamma(2\delta)\Gamma(\lambda-\delta)}{\Gamma(\lambda)\Gamma(\nu)\Gamma(\lambda+\delta+1)} \\ & \times {}_{q+3}F_{\mu+p+2} \left[\begin{matrix} \Delta(q; \gamma), \lambda+1, \lambda-\delta, 1; \\ \Delta(\mu; \nu), \Delta(p; \eta), \lambda, 1+\lambda+\delta; \end{matrix} \frac{q^q y}{\mu^\mu p^p a} \right], \end{aligned} \quad (3.2)$$

provided $\delta, \eta, \nu \in \mathbb{C}$ with $\min\{\Re(\delta), \Re(\eta)\} > 0$ and $a, \gamma, \lambda \in \mathbb{R}^+$ and $p, q, \mu \in \mathbb{N}$ with $\mu+p-q \geq 0$, $0 < \Re(\delta) < \lambda$, and $\Re(\nu) > -\mu$.

(iii) Setting $\nu = 0, \eta = \gamma = p = q = 1$ in (2.3) and using (1.9), we get

$$\begin{aligned} & \int_0^\infty x^{\delta-1} (x+a+\sqrt{x^2+2ax})^{-\lambda} E_\mu \left(\frac{y}{x+a+\sqrt{x^2+2ax}} \right) dx \\ &= 2^{1-\delta} a^{\delta-\lambda} \frac{\Gamma(2\delta)\Gamma(\lambda+1)\Gamma(\lambda-\delta)}{\Gamma(\lambda)\Gamma(1+\lambda+\delta)} \\ & \times {}_3F_{\mu+2} \left[\begin{matrix} \lambda+1, \lambda-\delta, 1; \\ \Delta(\mu; 1), \lambda, 1+\lambda+\delta; \end{matrix} \frac{y}{\mu^\mu a} \right], \end{aligned} \quad (3.3)$$

provided $\Re(\delta) > 0, \lambda \in \mathbb{R}^+$, and $\mu \in \mathbb{N}$.

In addition to three formulas (3.1), (3.2) and (3.3), by suitably specializing the integral formulas in Section 2, we can present other numerous integral formulas associated with extended Wright-Bessel functions and extended Mittag-Leffler functions.

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